

# Precise study of $^8\text{B}$ beta decay

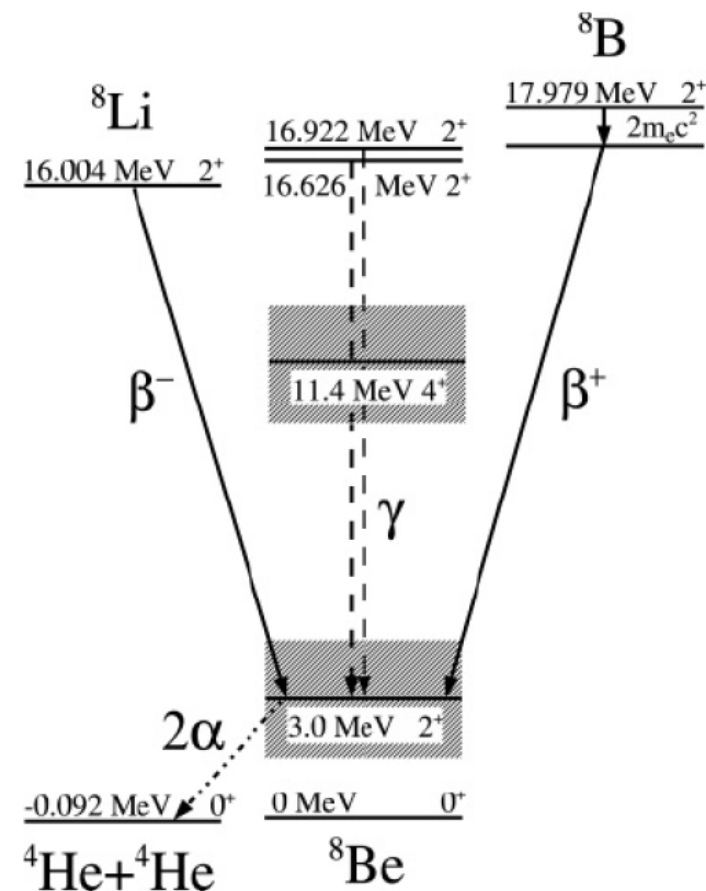
Thomas ROGER (K.U.Leuven)

- ${}^8\text{B}$  beta decay to  $2^+$  states in  ${}^8\text{Be}$
- Unable to reproduce beta-delayed  $2\alpha$  spectrum nor  $\alpha$ - $\alpha$  scattering cross section with 3 levels only

→ Add a 4<sup>th</sup> level

- at low excitation energy<sup>(1)</sup> : does not reproduce both  $\alpha$ - $\alpha$  scattering and beta-delayed  $\alpha$  spectra
- at higher excitation energy : physical meaning?

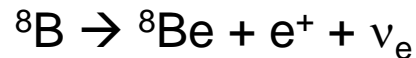
→ Direct decay via the continuum or higher lying resonances? <sup>(2)</sup>



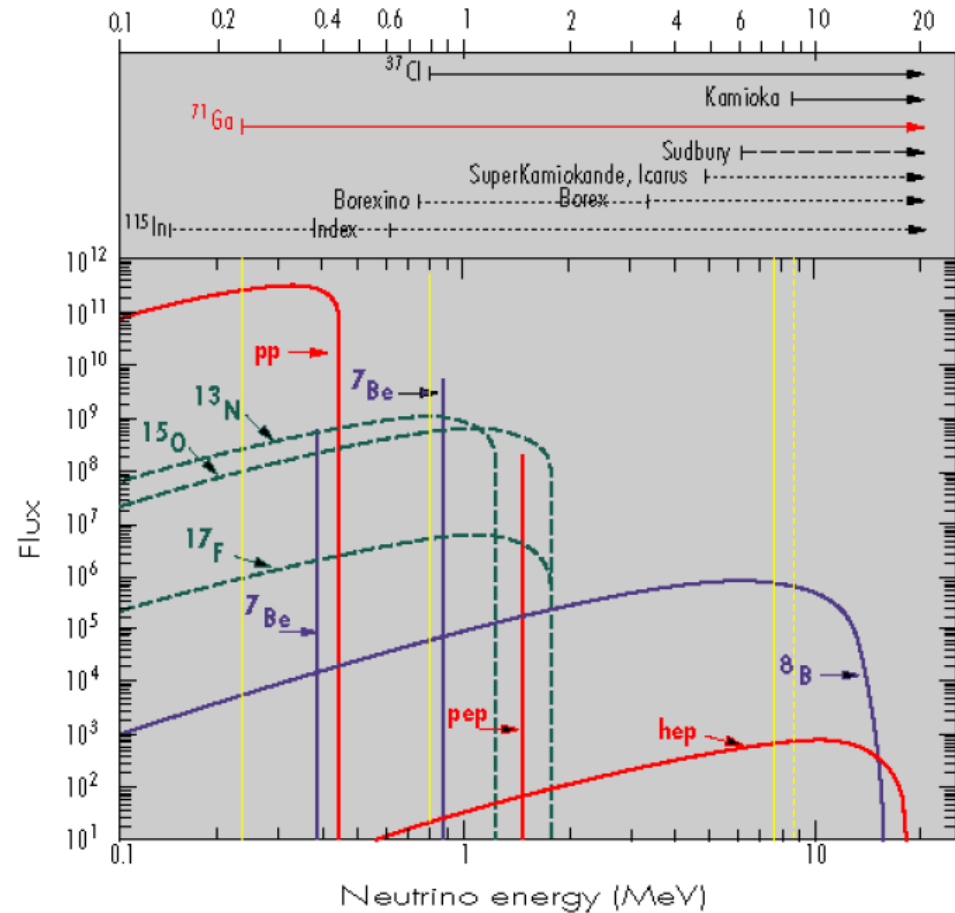
<sup>(1)</sup> F.C Barker, Aust. Journal of Physics **42**, 25 (1989)

<sup>(2)</sup> S. Hyldegaard, PhD thesis, Univ. Of Aarhus (2010)

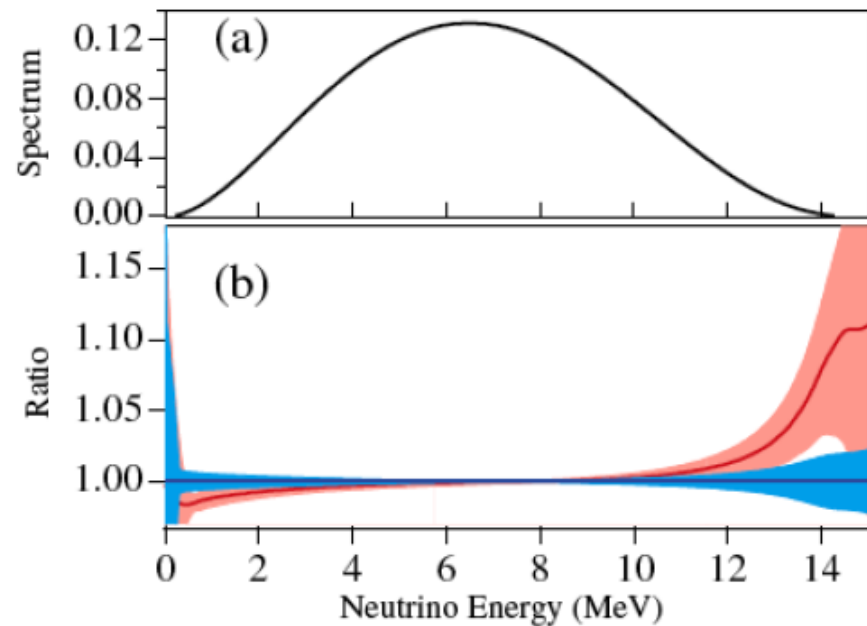
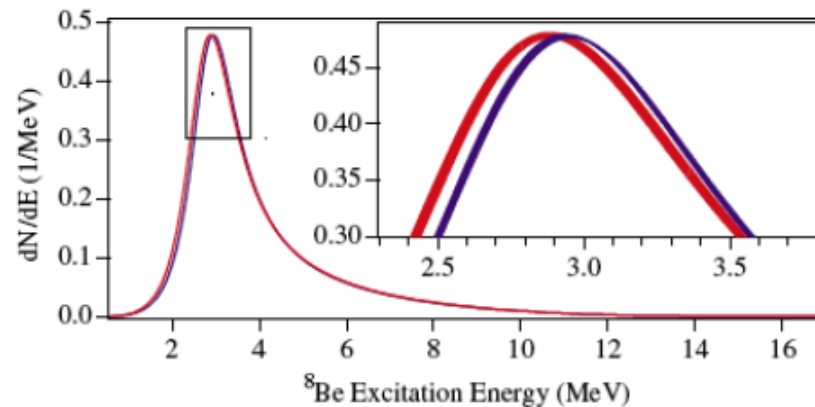
- Solar neutrino detectors  
Super-Kamiokande, SNO, ICARUS  
are sensitive to neutrinos from:



- The spectrum is depleted and distorted by neutrino oscillations. The distortion depends upon the oscillation parameters
- ➔ Need to know precisely the unperturbed  ${}^8\text{B}$   $\nu_e$  energy spectrum!



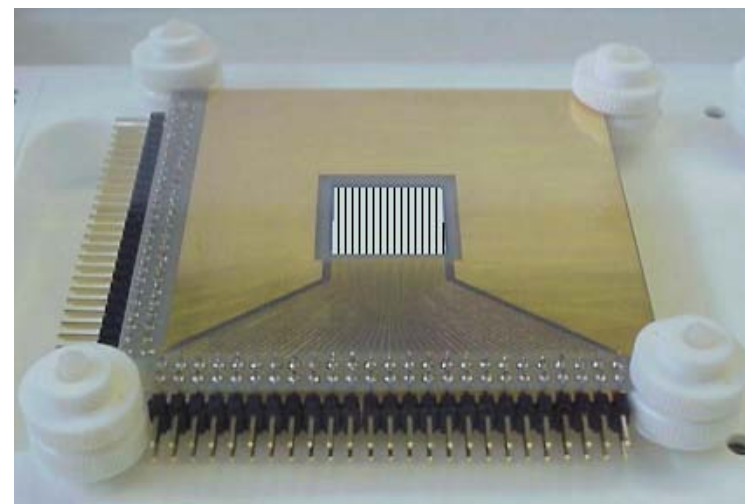
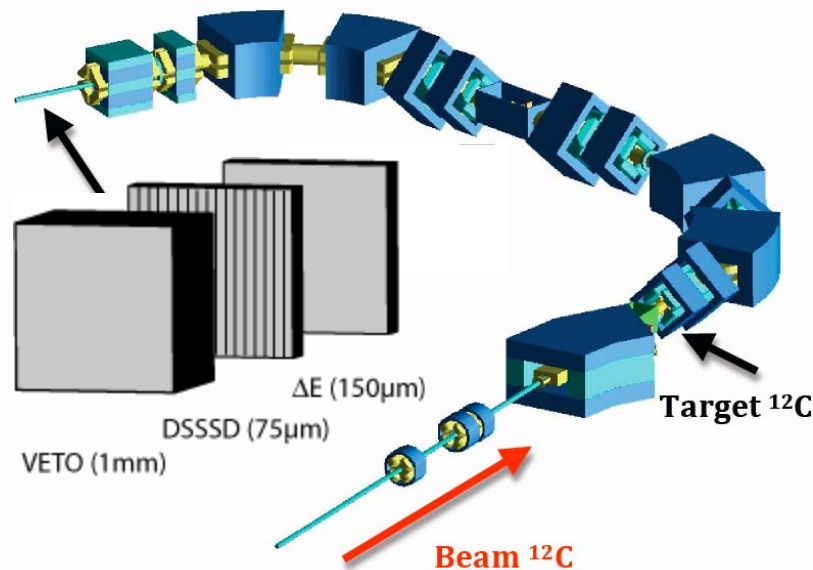
- 3 different measurements since 2000 :
  - Ortiz : coincident  $\alpha$  in C foil <sup>(1)</sup>
  - Winter : coincident  $\alpha$  in Si detector <sup>(2)</sup>
  - Bhattacharya : single  $\alpha$  in C foil<sup>(3)</sup>
- Large disagreement between Ortiz and Winter-Bhattacharya measurements
- All measurements have high detection threshold ( $>1\text{MeV}$ )  $\rightarrow$  large uncertainties in the high energy part of the  $\nu$  spectrum



- <sup>(1)</sup> C.E. Ortiz *et al.*, Phys. Rev. Lett. **85**, 2909 (2000)
- <sup>(2)</sup> W.T. Winter *et al.*, Phys. Rev. Lett. **91**, 252501 (2003)
- <sup>(3)</sup> M. Bhattacharya *et al.*, Phys. Rev. C **73**, 055802 (2006)

- Production of the  $^8\text{B}$  beam :  
fragmentation of  $^{12}\text{C}$  on  $^{12}\text{C}$
- Purification with the TRI $\mu$ P spectrometer <sup>(1)</sup>
- Identification via  $\Delta E$ -ToF &  $\Delta E$ -E
- Implantation in a DSSSD <sup>(2),(3)</sup>
  - 16x16 mm<sup>2</sup>, 78  $\mu\text{m}$  thick
  - 48x48 strips  $\rightarrow$  2304 pixels
  - Cooled electronics

Low detection threshold  
No dead layer effects



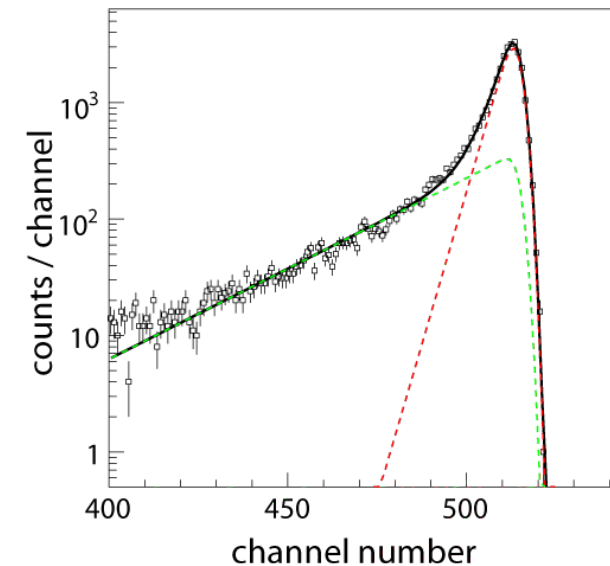
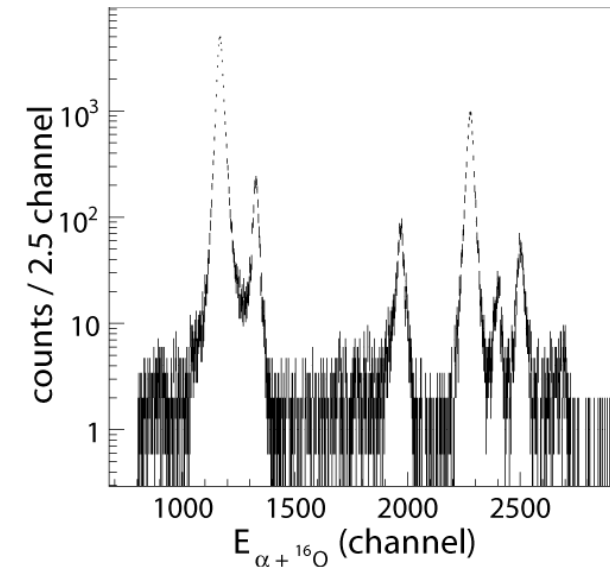
<sup>(1)</sup> G.P.A Berg. et al., NIM **A560**, 169 (2006)

<sup>(2)</sup> D. Smirnov et al., NIM **A547**, 480 (2005)

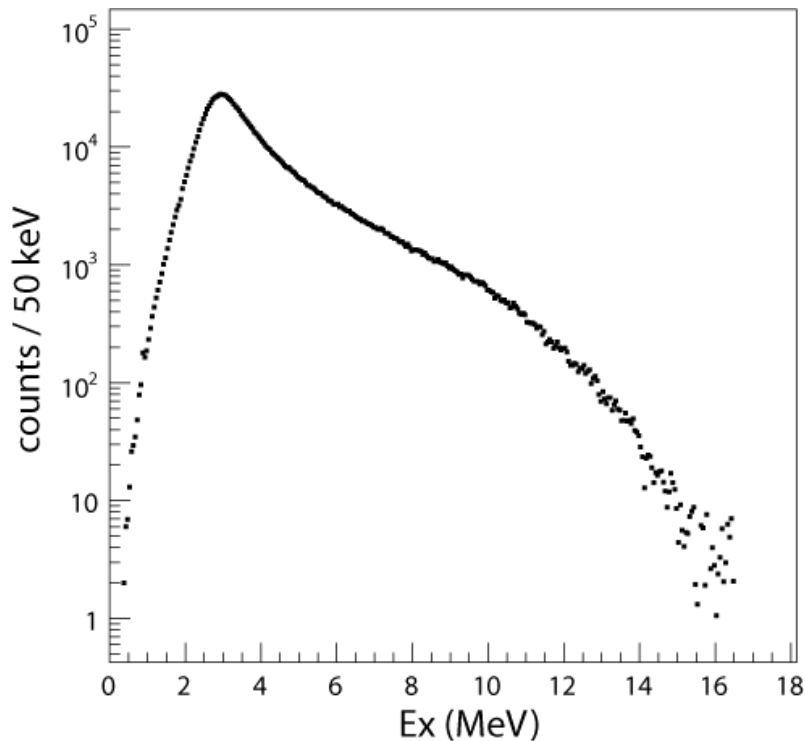
<sup>(3)</sup> J. Büscher et al., Nucl. Instr. and Meth. **B266**, 4652 (2008)

- Measurement using implanted  $^{20}\text{Na} \rightarrow ^{20}\text{Ne}^* \rightarrow ^{16}\text{O} + \alpha + e^+$ 
  - ➔  $^{16}\text{O}$  recoil pulse height defect:
    - ➔ TRIM:  $E_{\text{ionisation}} = 95\text{-}97\%$  ( $\sim 50$  keV)
  - ➔  $\beta$ - $\nu$  angular correlations &  $\beta$  summing:
    - ➔ Full G4 simulation ( $\sim 30$  keV)
    - ➔ Deconvolution
  
- Calibration with external  $\alpha$  sources
  - ➔ Tailing due to the dead layer ( $\sim 10$  keV)

Effects of the order of the precision needed



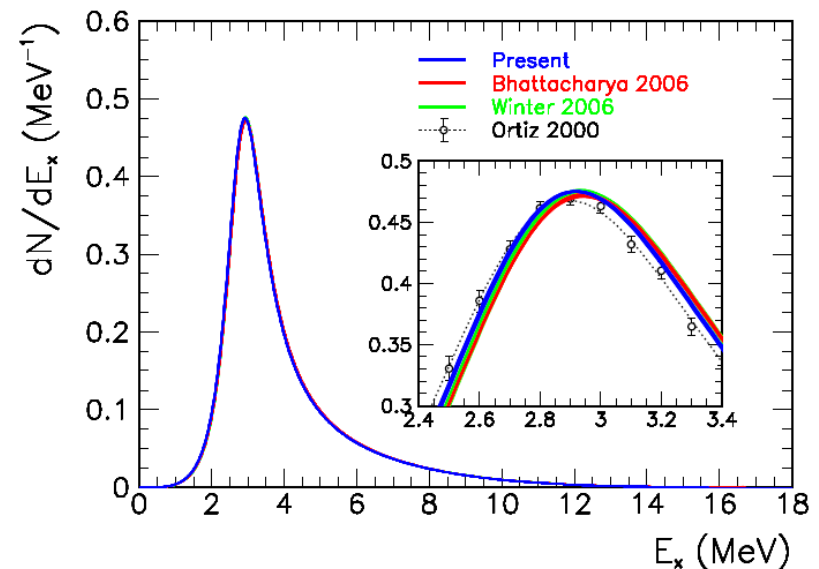
→  $10^6$  implantation-decay events



→ Detection threshold = 400 keV

R-matrix preliminary results:

- Comments to previous analysis:
  - 16 MeV doublet not well treated  
(Physical energy  $\neq$  R-matrix energy)
- Spectrum maximum lies in between Ortiz and Winter's measurements.



O. Kirsebom, PhD Thesis, Univ. Of Aarhus (2010)

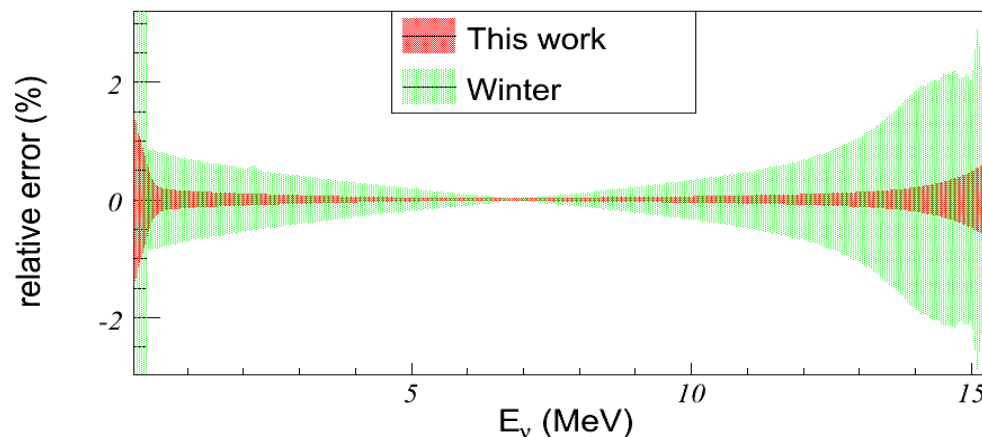
$$\left. \frac{dN}{dE_\beta} \right|_{Ex} \propto \underbrace{p_\beta E_\beta (E_0 - E_\beta)^2 \mathcal{F}(E_\beta)}_{\text{classic } \beta\text{-decay formula}} \times \underbrace{\left( 1 + \frac{3E - E_0 - 3\vec{p}_\beta \cdot \vec{k}_\nu}{M} \right)}_{\text{recoil energy}} \times \underbrace{C(E_0, E_\beta)}_{\text{recoil order correction}} \times \underbrace{\mathcal{R}(E_0, E_\beta)}_{\text{radiative correction}}$$

Sources of error:

- Theoretical: error on recoil order correction dominates
- Calibration error:  $\sim 7$  keV @ 3 MeV
- Error at high  $E_\nu$  (lack of experimental data):

➔ Previously: changing channel radius in the R-matrix fit

➔ We extract the  $\nu$  spectrum directly from the data





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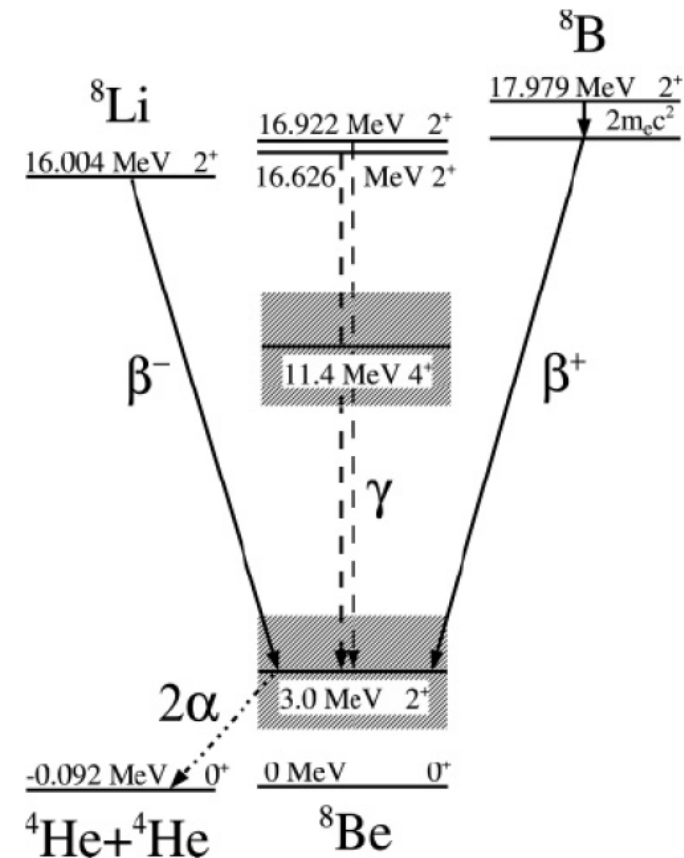
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- Fit of the distribution in the R-matrix framework
  - ➔ One channel – Many level approach
  - ➔  $\beta$  decay selection rules :  
decay to  $0^+$  &  $4^+$  are secondly forbidden
  - ➔ Only  $2^+$  states are energetically allowed
  - ➔ 3 physical states + 1 background state



$$\frac{dN}{dE_x} = \left( \frac{Nt_{1/2}}{6166} \right) f_{\beta}(E_x)(a^2(E_x) + c^2(E_x))$$

Fermi matrix element

$$a^2(E_x) = \frac{P(E_x)}{\pi} \times \left( \frac{\left| \sum_{j=1}^4 \frac{M_{Fj} \gamma_j}{\varepsilon_j - E_x} \right|^2}{\left| 1 - (S(E_x) - B + iP(E_x)) \sum_{j=1}^4 \frac{\gamma_j^2}{\varepsilon_j - E_x} \right|^2} \right)$$

GT matrix element

$$c^2(E_x) = \frac{P(E_x)}{\pi} \times \left( \frac{\left| \sum_{j=1}^4 \frac{M_{GTj} \gamma_j}{\varepsilon_j - E_x} \right|^2}{\left| 1 - (S(E_x) - B + iP(E_x)) \sum_{j=1}^4 \frac{\gamma_j^2}{\varepsilon_j - E_x} \right|^2} \right)$$

Penetration factor

$$P(E_x) = \sqrt{\frac{2M_{2\alpha}E_x}{\hbar^2}} \frac{r_c}{F(r_c)^2 + G(r_c)^2}$$

Shift factor

$$S(E_x) = \sqrt{\frac{2M_{2\alpha}E_x}{\hbar^2}} \frac{r_c [F(r_c)F'(r_c) + G(r_c)G'(r_c)]}{F(r_c)^2 + G(r_c)^2}$$

States parametrisation - 4x4 parameters:

$2^+_1$

- $\varepsilon_1$  free
- $\gamma_1^2$  free
- $\mathcal{M}_{GT1}$  free
- $\mathcal{M}_{F1} = 0$  (SM.consideration)

$2^+_4$

- $\varepsilon_4 = 37.0$  (syst. – weak influence)
- $\gamma_4^2$  free
- $\mathcal{M}_{GT4}$  free
- $\mathcal{M}_{F4} = 0$  (hypothesis)

$2^+_2 - 2^+_3$  : max. mixed isospin doublet

- |  |  |
|--|--|
| - $\varepsilon_2 = 16.626$ (exp)                 | - $\varepsilon_3 = 16.922$ (exp)                 |
| - $\gamma_2^2 = 10.96$ keV (exp)                 | - $\gamma_3^2 = 7.42$ keV (exp)                  |
| - $\mathcal{M}_{GT2} = \alpha \mathcal{M}_{GTA}$ | - $\mathcal{M}_{GT3} = \beta \mathcal{M}_{GTA}$  |
| - $\mathcal{M}_{F2} = \beta \mathcal{M}_{GTB}$   | - $\mathcal{M}_{F3} = -\alpha \mathcal{M}_{GTB}$ |

With mixing parameters :  $\alpha^2 = \Gamma_2/(\Gamma_2+\Gamma_3)$  and  $\beta^2 = \Gamma_3/(\Gamma_2+\Gamma_3)$

and  $\left\{ \begin{array}{l} \mathcal{M}_{GTB} = \sqrt{2} \text{ (T=1 isospin eigenstate is the IAS of } ^8\text{Li \& } ^8\text{B GS)} \\ \mathcal{M}_{GTA} \text{ free} \end{array} \right.$